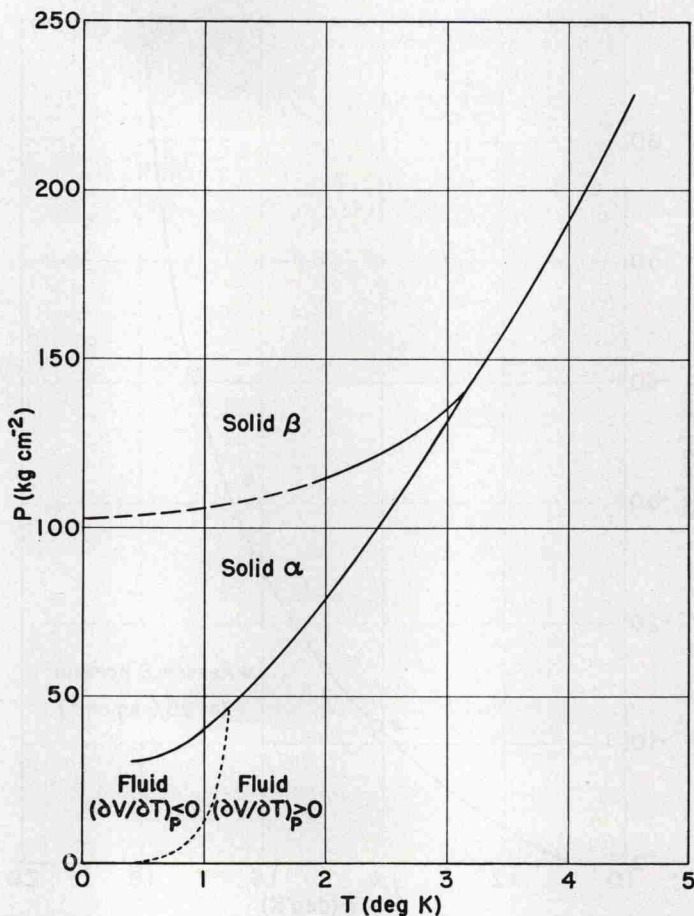


FIG. 8.  $\alpha_f$  as a function of  $T$  for He<sup>4</sup> at  $P = 25.47 \text{ kg/cm}^2$  and  $T < T_\lambda$ .

perature. The curve of Fig. 3 exhibits a sharp rise in  $\alpha_f$  with decreasing  $P_m$  until a maximum is reached at  $\sim 50 \text{ kg/cm}^2$ . At pressures below this,  $\alpha_f$  decreases rapidly, becomes discontinuous at the  $\lambda$ -point by assuming large negative values which decrease in magnitude with further decreases in  $P_m$ , as shown in Fig. 8. The  $\alpha_f$  values of Keesom and Keesom (9), plotted in Fig. 8, were derived from their PVT measurements and show good agreement with the present data.

The  $V_f$  measurements of Table I agree with those of Dugdale and Simon (3) to within 1 percent. The newly determined melting curve, expressed by Eq. (2) and Table VI, forms a smooth extension of that reported earlier (1).

FIG. 9. Condensed phase diagram for He<sup>3</sup>.

## B. DISCUSSION OF He<sup>3</sup> RESULTS

At high pressures the properties of He<sup>3</sup> along the melting curve roughly parallel those of He<sup>4</sup>. At low pressures, although He<sup>3</sup> does not display superfluidity, there are other unique features which merit detailed discussions.

### 1. The solid-solid transition

With reference to Fig. 2, the discontinuity in the  $\Delta V_m$  curve for He<sup>3</sup> at  $\sim 141$  kg/cm<sup>2</sup> is a consequence of a triple point in the melting curve where two types of solid are in equilibrium with the fluid phase. A careful determination of the melting curve showed a slight discontinuity in slope which occurs at 141 kg/cm<sup>2</sup>